# An area law and sub-exponential algorithm for 1D systems. 

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And yet . . . condensed matter physicists do it!
Heuristic techniques, DMRG, have been very successful for 1D systems.

Is there a principled phenomenon behind this?
Is there a clean well defined class of quantum many body systems that we can analyze?

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A first point of entry is a remarkable conjecture:

Area Law: Given a gapped local Hamiltonian, for any subset $S$ of particles, the entanglement entropy of $\rho_{S}$, the reduced density matrix of the ground state restricted to $S$, is bounded by the surface area of $S$ i.e. the number of local interactions between $S$ and $\bar{S}$.

## Basic Questions



- Can you prove an area law?
- If so, do these states have small working descriptions?
- Can they be efficiently computed?


## Concretely in 1D



Given:

- $n d$-dimensional particles on a line, $\mathcal{H}=\left(\mathbb{C}^{d}\right)^{\otimes n}$,
- local operators $0 \leq H_{i} \leq 1$ acting non-trivially on the $i$ th and $i+1$ st particle.
- a Hamiltonian $H=\sum_{i} H_{i}$ with a gap $\epsilon$ between the energy of the ground state and the next lowest energy.


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- a Hamiltonian $H=\sum_{i} H_{i}$ with a gap $\epsilon$ between the energy of the ground state and the next lowest energy.
Goal: structural properties of the ground state $\mid \Gamma>$.


## Result: 1D Area Law

Previous results for 1D:

- Hastings (2007) with bound $e^{O(\log d / \epsilon)}$.
- Existence of an MPS with polynomial bond dimension.
- Finding an approximation to the ground state is $\in N P$.
- Arad, Landau, Vazirani (2011): $\tilde{O}\left(\frac{\log d}{\epsilon}\right)^{3}$ for frustration free system.
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This result:

Theorem: The entanglement entropy of the ground state of a 1D gapped Hamiltonian is bounded by $\tilde{O}\left(\frac{\log ^{3} d}{\epsilon}\right)$

- Exponential improvement of the bound.
- Bound the cusp of a 2D sub-volume law.
- Implies a sublinear bond dimension MPS which leads to . . .


## Sub-exponential algorithm

Theorem: There is a subexponential time algorithm for finding an inverse polynomial approximation to ground state of a 1D gapped Hamiltonian.

Combines sublinear bond dimension with a dynamical programing algorithm (Aharonov, Arad, Irani, 2009).

## Preliminaries: Entanglement Rank

For a vector $v \in \mathcal{H}_{1} \otimes \mathcal{H}_{2}$ with Schmidt decomposition $v=\sum_{i=1}^{D} a_{i} \otimes b_{i}$, has entanglement rank $D$. Operators

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- Operators of the form $\sum_{1}^{C} A_{i} \otimes B_{i}$ can only increase the entanglement rank by a factor of $C$.

- Local operator $H_{i}$ can only increase the entanglement rank across $i, i+1$ by $d^{2}$.



## Preliminaries: Functional calculus of an operator

What does the operator $f(H)=H^{2}-2 H-1$ look like?

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Depiction of $\mathrm{f}(\mathrm{H})$

Eigenspaces of H

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Depiction of $\mathrm{f}(\mathrm{H})$

(Located at corresponding eigenvalues)

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We are looking for an operator $K$ with 2 properties:

- It approximately projects onto the ground state:

- It doesn't increase the entanglement too much:


Such an operator is a $(D, \Delta)$ Approximate Ground State Projection (AGSP).

## The consequence of a good AGSP: An area law

Theorem (Area Law) [Arad, Landau, Vazirani] The existence of an AGSP $K$ for which $D \Delta<1 / 2$ proves that the ground state has entropy $O(1) \log D$.

## Building good AGSP's: reduce the norm

Looking for low entanglement operators that look like:


Smaller $\|H\|$ would be better but we don't want to lost the local structure around the cut.
Solution: Replace $H=\sum_{i} H_{i}$ with $H^{\prime}=H_{L}+H_{1}+H_{2}+\cdots+H_{s}+H_{R}$.


## Building good AGSP's: Chebyshev polynomials

Chebyshev polynomials: small in an interval:
Chebyshev Polynomials


## A good AGSP

A dilation and translation of the Chebyshev polynomial gives:

with

$$
\Delta=e^{-\frac{\ell \sqrt{\epsilon}}{\sqrt{\left\|H^{\prime}\right\|}}} .
$$

## Entanglement Increase due to a single term of $\left(H^{\prime}\right)^{\ell}$

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\left(H^{\prime}\right)^{\ell}=\sum\left(\text { product of } H_{j}\right) .
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For a single term:


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Total: $d^{2 \ell / s+s}$


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Problem: Too many $\left(s^{\ell}\right)$ terms in naive expansion of $\left(H^{\prime}\right)^{\ell}$.
Need to group terms in a nice way but it all works out with total entanglement increase of the same order as the single term.

## Putting things together: Area Law for $H^{\prime}$

Chebyshev $C_{\ell}\left(H^{\prime}\right)$ has $\Delta \approx e^{-O(\ell \sqrt{\epsilon} / \sqrt{s})}$ :


Entanglement analysis yields $D \approx O\left(d^{\ell / s+s}\right)$.


Chosing $\ell=s^{2}$ yields $\log (D \Delta) \approx-s^{3 / 2} \sqrt{\epsilon}+s \log d$. Approximate equality occurs with $s \approx \log ^{2} d / \epsilon$ which yields $D \approx \log ^{3} d / \epsilon$.

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Definition of $H_{L}$ and $H_{R}$ using truncation:


$$
H_{L}=\left(\sum_{i<1} H_{i}\right)^{\leq t}, \quad H_{R}=\left(\sum_{i>s+1} H_{i}\right)^{\leq t} .
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Question How does the Hamiltonian $H^{\prime}=H_{L}+H_{1}+\cdots+H_{s}+H_{R}$ compare to $H=\sum_{j} H_{j}$ ?

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Robustness Theorem: The gaps of $H$ and $H^{\prime}$ are of the same order and the ground states of $H$ and $H^{\prime}$ are within $\exp (-t)$.

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Robustness Theorem: The gaps of $H$ and $H^{\prime}$ are of the same order and the ground states of $H$ and $H^{\prime}$ are within $\exp (-t)$.

Area law for $H$ now follows by starting with a constant truncation level $t=t_{0}$ and then letting it grow to $t=O(\log n)$.

## Summary



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Area law


Subexponential time algorithm

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Schuch, Cirac, Verstraete

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Good AGSP


The structural engine for these results are AGSP's.

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- Of independent interest: entanglement rank has a "random walk" type behavior (added entanglement of $H^{\ell}$ is $d^{O(\sqrt{\ell})}$ ).
- Of independent interest: robustness theorem of truncation.

